

X-Ray Photoelectron Studies of Platinum and Palladium Complexes; Observation of the *trans*-Influence and Distinction Between Terminal and Bridging Chlorine

By D. T. CLARK* and D. B. ADAMS

(Department of Chemistry, University of Durham, South Road, Durham)

and D. BRIGGS

(Imperial Chemical Industries, Ltd., Petrochemical and Polymers Laboratory, Runcorn Heath, Cheshire)

Summary Molecular core binding energies have been measured for a series of platinum and palladium complexes; the results show the great potential of X-ray photoelectron spectroscopy for studies in structure and bonding.

THE advantage of X-ray photoelectron spectroscopy^{1,2} is that in principle it is capable of revealing something of the electronic environment of each constituent atom of a molecule and this makes a study of structure and bonding in Pt^{II} and Pd^{II} complexes particularly attractive. The relatively small amounts of sample required coupled with the large ionization cross sections for the Pt(4f) and Pd(3d) core levels are then additional bonuses. We report here preliminary results concerning bonding between Pt and alkenes and alkynes, the *trans*-influence in square-planar Pt^{II} and Pd^{II} complexes, and distinction between terminal and bridging chlorine in *trans*-L₂M₂Cl₄ (L = R₃P or Me-CHCH₂; M = Pt or Pd).

Spectra were measured on an A.E.I. ES 100 electron spectrometer for samples studied as thin films on a conducting backing using MgK $\alpha_{1,2}$ radiation. The sensitivity of the technique is illustrated by the spectra shown in the Figure for the Pt(4f) and I(3d) core levels of *trans*-(PMe₃)₂-PtI₂ where the sample size was *ca.* 1 mg.

The results are shown in the Table. It may be shown² that the binding energy of a given core level of an atom in a molecule is related to the charge distribution by equation (1) where E_0^i is a reference level and the second term

$$E^i = E_0^i + k_{qi} + \sum_{i \neq j} \frac{q_j}{r_{ij}} \quad (1)$$

represents the potential from the charge at the atom considered. The third term accounts for the potential from the rest of the molecule (an intramolecular Madelung type potential; for ionic solids however the summation should in principle be extended to all the atoms in the crystal lattice).

In general the *shift* in binding energy of a given core level is dominated by the charge at the atom concerned, the potential from the other atoms being usually much smaller

also investigated some of these compounds, *e.g.* *cis* (Ph₃P)₂PtL₂, L₂ = PhC≡CPh, CH₂=CH₂, and Cl₂. In general the agreement between our results is satisfactory with the

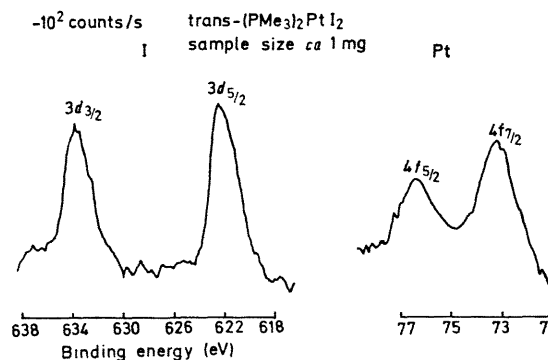
		Binding energies (in eV)				
		L ₂	Pt(4f _{7/2})	C(1s)	P(2p _{3/2})	L
<i>cis</i> -(Ph ₃ P) ₂ PtL ₂		H-C≡C-H	72.3	284.7	131.0	C(1s) C≡C 283.3*
		CH ₃ -C≡C-H	72.4	284.7	130.9	C(1s) C≡C 283.3*
		Ph-C≡C-H	72.4	284.7	130.9	C(1s) C≡C 283.3*
		H ₂ C=CH ₂	72.4	284.7	131.1	C(1s) C=C 283.2*
		(Cl) ₂	72.2	284.7	131.0	Cl(2p _{3/2}) 198.0
		<i>trans</i> -(Me) ₂	71.2	284.7	130.7	
		(Cl) ₂	72.1	284.7	130.9	Cl(2p _{3/2}) 198.0
		(Me) ₂	71.2	284.7	130.5	
(Ph ₂ P)-C ₂ H ₄ -(PPh ₂)						
L ₂ PtX ₂						
L	X	Pt(4f _{7/2})	Cl(2p _{3/2})	P(2p _{3/2})	L	C(1s)
<i>cis</i> -Bu ⁿ ₃ P	Cl	72.0	198.1	131.1		284.7†
<i>trans</i> -Bu ⁿ ₃ P	Cl	72.0	199.3	130.7		284.7†
		Pt(4f _{7/2})	I(3d _{5/2})			
<i>cis</i> -Me ₃ P	I	72.9	621.4	132.6		285.1
<i>trans</i> -Me ₃ P	I	72.9	622.1	132.1		285.0
	Pt metal	71.1				
	(PhEt ₂ P) ₂ PdCl ₂	Pd(3d _{5/2})	Cl(2p _{3/2})	P(2p _{3/2})		C(1s)
<i>cis</i>		338.0	197.4	130.9		284.7†
<i>trans</i>		338.0	198.2	130.6		284.7†
<i>trans</i> -L ₂ M ₂ Cl ₄						N.q.r. frequency
M = Pt	Pt(4f _{7/2})	Cl(2p _{3/2})		L		³⁵ Cl ν ₀ MHz
L = MeCH=CH ₂	72.7	199.4 terminal	C(1s) CH ₃	285.0	C=C	24.12 terminal ^{5,8}
		198.4 bridging				15.95 bridging ^{5,8}
L = Pr ⁿ ₃ P	71.9	199.1 terminal	C(1s)	284.7†	P(2p _{3/2})	22.36 terminal ^{5,8}
		198.1 bridging				15.46 bridging
M = Pd	Pd(3d _{5/2})	Cl(2p _{3/2})		L		
L = MeCH=CH ₂	337.9	199.3 terminal	C(1s) CH ₃	285.0	C=C	284.2
		198.3 bridging				
L = Bu ⁿ ₃ P	337.7	198.7 terminal	C(1s)	284.7†	P(2p _{3/2})	130.7
		197.7 bridging				19.50 terminal ^{5,8}
						12.77 bridging ^{5,8}
Pd metal	336.3					

* These appear as shoulders on the main peak due to the Ph carbons. Deconvolutions were carried out with a Dupont 310 curve resolver. All levels were calibrated with respect to Au(4f_{7/2}) 84 eV binding energy.

† Centroid of relatively broad peak half width 1.8 eV.

and opposite in sign. In the closely related series of molecules described here, therefore, the shifts in a given core level will qualitatively reflect differences in charge distribution. With large polyatomic ligands such as (Ph₃P) charge migration to or from the ligand has little effect on any individual carbon atom, since it is spread over so many centres, and hence within experimental error the C(1s) binding energy for the phenyl carbons is identical for all samples and consists of a broad peak (half width *ca.* 1.8 eV) centred at 284.7 eV binding energy. Comparison of the core binding energies for *trans*-(Ph₃P)₂PtMe₂ and the related chelated phosphine derivative which of necessity has *cis*-stereochemistry indicates that the core levels of the metal are not greatly influenced by the stereochemistry and are comparable in binding energy to the metal itself. On replacement of the strong σ donor CH₃ ligands† by an alkene or alkyne there is considerable electron transfer from metal to ligand. This is shown by the increase in binding energy for the metal core levels and the low binding energy for the alkene and alkyne carbon atoms. [For comparison the C(1s) binding energies for -C≡C- and >C=C< in the free ligands would be *ca.* 284.9 eV]. This result may be understood in terms of the Dewar-Chat-Duncanson model³ in which back donation from filled Pt *d* orbitals into the π* antibonding orbitals of the hydrocarbon is the dominating feature. Siegbahn and his co-workers⁴ have

exception of *cis*-(Ph₃P)₂PtCl₂ for which a Pt(4f_{7/2}) binding energy of 73.4 eV was obtained. This discrepancy is



FIGURE

difficult to understand since we have investigated numerous compounds of general formula (PR₃)₂PtCl₂ (both *cis* and *trans*) and in each case the binding energies for the Pt(4f_{7/2}) levels are close to 72.2 eV.

X-Ray crystallographic data show quite clearly that there is a considerable increase in metal-halogen bond

† The strong σ-donation from the methyl groups would undoubtedly be reflected in a higher C(1s) binding energy for these ligands. Unfortunately this is obscured by the much higher intensity peak due to the phenyl carbons.

length in going from the *trans* to *cis* isomer in square-planar Pt^{II} and Pd^{II} complexes of the type (PR₃)₂MX₂. For X = Cl, n.q.r. data suggest⁵ that the bond lengthening is accompanied by a considerable electron drift to chlorine. The data in the Table show that in a given *cis-trans* pair the binding energy of the metal is the same but that the binding energy of the halogen is considerably lower in the *cis*-isomer. By contrast the binding energy for the phosphorus atom of the phosphine ligands is somewhat higher in the *trans*-isomers; however, the absolute values for the shift in binding energy within a given *cis-trans* pair is smaller than for the halogen. This undoubtedly arises from the dispersal of any charge over the alkyl and aryl groups of the phosphine ligands. Considering first (R₃P)₂MCl₂ (M = Pt or Pd), the results complement and extend the picture of charge distribution in a given *cis-trans* pair suggested by n.q.r. which provides information on only the chlorine atoms. The Pt and Pd core binding energies indicate that the electronic environment about the metal is the same for a given *cis-trans* pair and that an increase in density around chlorine in going from the *trans* to *cis* isomer (as shown by lower binding energies) is compensated by a decrease in density around phosphorus. This is readily understood in terms of the strong σ donor properties of phosphine ligands and a Syркиn type rehybridization scheme.⁶ Comparing corresponding isomers[†] (R₃P)₂MCl₂ (M = Pt and Pd) for the Pt(4f_{7/2}) and Pd(3d_{5/2}) core levels the shift with respect to the metal is 0.9 and 1.7 eV, respectively, indicating a positive charge, that on palladium being the larger. By contrast the Cl(2p_{3/2}) binding energies are lower for the palladium compounds indicative of a higher electron density compared with the corresponding platinum compound.

For (R₃P)₂PtX₂ (X = Cl or I) the binding energies suggest greater positive charge on the platinum and phosphorus atoms for the iodide. This can only reasonably be accommodated by a relatively greater electron density on the halogen atoms in the iodides. This confirms Parshall's⁷ order of σ donation of Cl⁻ > I⁻ (the π acceptor properties being much smaller) obtained from ¹⁹F n.m.r. studies.

The measured molecular core binding energies for *trans*-L₂M₂Cl₄ (L = R₃P or MeCHCH₂; M = Pt or Pd) are also shown in the Table. In square-planar Pt^{II} and Pd^{II} chlorides, the Cl(2p) core levels appear as spin-orbit-split (*ca.* 1.8 eV) doublets. In the bridged species, however, the Cl(2p) levels appear as broad unresolved peaks, which may be deconvoluted to give four components, corresponding to the spin-orbit-split components for two non-equivalent

chlorines. In each case the difference in binding energy between terminal and bridging chlorines is *ca.* 1.0 eV.

On the basis of n.q.r. data^{5,8} (also shown in the Table) the chlorine with lower binding energy may be assigned as bridging in each case. For the phosphine complexes, the binding energies of comparable terminal and bridging chlorines is somewhat lower for the palladium compound. This follows the behaviour for the square-planar complexes [(R₃P)₂MCl₂]. This reflects the higher electron density around chlorine and is also evident in the n.q.r. results.^{5,8} The binding energy of the metal core levels (4f_{7/2}) and Pd(3d_{5/2}) suggests that in the chlorine-bridged phosphine complexes the electronic environment around the metal is similar to that in the corresponding square-planar complexes.

For the propene complexes, strong interaction between filled metal *d* orbitals and antibonding π^* orbitals is shown by the relatively low C(1s) binding energy of the C=C carbons although comparison with (Ph₃P)₂PtC₂H₄ suggests that the interaction is smaller in the bridged complexes. Electron migration to the alkene ligand in the platinum complex is largely at the expense of the metal, since the Pt(4f_{7/2}) level binding energy is considerably increased compared with the phosphine complex, the Cl(2p_{3/2}) levels being shifted to higher binding energy by a small amount. In the corresponding palladium compound the situation is reversed and electron migration to the alkene is at the expense of chlorine, the binding energy of the palladium being similar in the two compounds. The explanation for this difference in behaviour can be understood by a consideration of the binding energy of the metal core levels with respect to the elements. For the phosphine complexes the shift to higher binding energy amounts to 0.8 and 1.3 eV for Pt(4f_{7/2}) and Pd(3d_{5/2}) respectively, indicative of considerable positive charge, that on palladium being the greater. Replacement of a strong σ donor ligand by a strong π acceptor ligand in going from the phosphine to alkene complex will therefore be expected to make an increased electron demand on the metal and other ligands. The latter will be much larger for the palladium compound since the metal carries a much larger positive charge than in the corresponding platinum complex.

We thank the S.R.C. for provision of equipment and for a research studentship (to D.B.A.). Thanks are due to Dr. P. L. Goggin, University of Bristol for providing samples of *cis*- and *trans*-(PMe₃)₂PtI₂ and to Professor K. Siegbahn for a preprint of his work.

(Received, March 1st, 1971; Com. 161.)

† We are of course comparing compounds with different phosphine ligands, the electronic effects of which will also be different. However these differences will be small as judged by the metal core bonding energies [*cf.* binding energies of *cis*-(Ph₃P)₂PtCl₂].

¹ K. Siegbahn, C. Nordling, A. Fahlman, R. Nordberg, K. Hamrin, J. Hedman, G. Johansson, T. Bergmark, S. E. Karlsson, I. Lindgren, and B. Lindberg, "ESCA; Atomic Molecular and Solid State Structure Studied by means of Electron Spectroscopy". Almqvist and Wiksells, Uppsala, 1967.

² K. Siegbahn, C. Nordling, G. Johansson, J. Hedman, P. F. Heden, K. Hamrin, U. Gelius, T. Bergmark, L. O. Werme, R. Manne, Y. Baer, "ESCA Applied to Free Molecules," North Holland, Amsterdam, 1969.

³ M. J. S. Dewar, *Bull. Soc. chim. France*, 1951, **18C**, 79; J. Chatt and L. A. Duncanson, *J. Chem. Soc.*, 1953, 2939.

⁴ C. D. Cook, K. Y. Wan, U. Gelius, K. Hamrin, G. Johansson, E. Olson, H. Siegbahn, C. Nordling, and K. Siegbahn, *J. Amer. Chem. Soc.*, 1971, submitted for publication.

⁵ C. W. Fryer and J. A. S. Smith, *J. Chem. Soc. (A)*, 1970, 1030.

⁶ (a) Y. K. Syркиn, *Bull. Acad. Sci., U.S.S.R., Classe Sci. chim.*, 1948, 69; (b) L. M. Venanzi, *Chem. in Britain*, 1968, **4**, 162.

⁷ (a) G. W. Parshall, *J. Amer. Chem. Soc.*, 1964, **86**, 5367; (b) R. V. Lindsay, jun., G. W. Parshall, and U. G. Stolberg, *J. Amer. Chem. Soc.*, 1965, **87**, 658.

⁸ C. W. Fryer, unpublished data.